

# Lecture 8: Diatomic Partition Functions

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Rotational and Vibrational Partition Functions

# Diatomic Ideal Gas

Diatomic Ideal gases have all degrees of freedom:

$$q = q_{\text{trans}} \cdot q_{\text{rot}} \cdot q_{\text{vibr}} \cdot q_{\text{elec}}$$

$$\varepsilon = \varepsilon_{\text{trans}} + \varepsilon_{\text{rot}} + \varepsilon_{\text{vibr}} + \varepsilon_{\text{elec}}$$

Translational and Electronic partition functions behave similarly to monoatomics

$$q_{\text{trans}} = \left[ \int_0^\infty \exp\left(-\beta \frac{h^2}{8ma^2} n^2\right) dn \right]^3$$

$$\langle \varepsilon_{\text{trans}} \rangle = \frac{3}{2\beta} = \frac{3}{2} k_B T \quad \frac{1}{2} k_B T \text{ per DOF}$$

$$q_{\text{elec}} = \sum_{i, \text{levels}} g_i e^{-\beta \varepsilon_i}$$

Virtually no contribution to  $\varepsilon$

# Vibrational Partition Function

Harmonic Oscillator:

$$\varepsilon_n = h\nu \left( n + \frac{1}{2} \right)$$

Substitute into partition function:

$$\begin{aligned} q_{\text{vib}} &= \sum_{n=0}^{\infty} e^{-\beta h\nu \left( n + \frac{1}{2} \right)} \\ &= \sum_{n=0}^{\infty} e^{-\beta h\nu n} \cdot e^{-\beta h\nu \frac{1}{2}} \end{aligned}$$

Factor out constant  $e^{-\beta h\nu/2}$ :

$$\begin{aligned} &= e^{-\beta h\nu/2} \sum_{n=0}^{\infty} e^{-\beta h\nu n} \\ &= e^{-\beta h\nu/2} \sum_{n=0}^{\infty} \left( e^{-\beta h\nu} \right)^n \end{aligned}$$

# Vibrational Partition Function

Vibrational partition function:

$$q_{\text{vib}} = e^{-\beta h\nu/2} \sum_{n=0}^{\infty} \left( e^{-\beta h\nu} \right)^n$$

Recognize as a geometric series:

$$\sum_{n=0}^{\infty} r^n = \frac{1}{1-r} \quad |r| < 1$$

Where  $r = e^{-\beta h\nu}$ :

$$\sum_{n=0}^{\infty} \left( e^{-\beta h\nu} \right)^n = \frac{1}{1 - e^{-\beta h\nu}}$$

Thus the vibrational partition function becomes:

$$q_{\text{vib}} = \frac{e^{-\beta h\nu/2}}{1 - e^{-\beta h\nu}}$$

## What is Vibrational Temperature?

Define the vibrational temperature  $\Theta_{\text{vib}} = \frac{h\nu}{k_B}$ :

$$q_{\text{vib}} = \frac{e^{-\frac{\Theta_{\text{vib}}}{2T}}}{1 - e^{-\frac{\Theta_{\text{vib}}}{T}}}$$

Consider the probability of finding a particle in an excited state:

$$P_{\text{excited}} = 1 - P_0$$

Ground State energy:

$$\varepsilon_0 = h\nu \left( 0 + \frac{1}{2} \right) = \frac{h\nu}{2}$$

Probability of being in the ground state:

$$P_0 = \frac{e^{-\beta\varepsilon_0}}{q_{\text{vib}}} = \frac{e^{-\frac{\beta h\nu}{2}}}{\left( \frac{e^{-\frac{\Theta_{\text{vib}}}{2T}}}{1 - e^{-\frac{\Theta_{\text{vib}}}{T}}} \right)}$$

# What is Vibrational Temperature?

Substitute vibrational temperature  $\Theta_{\text{vib}} = \frac{h\nu}{k_B}$ :

$$P_0 = \frac{e^{-\frac{\Theta_{\text{vib}}}{2T}}}{\left( \frac{e^{-\frac{\Theta_{\text{vib}}}{2T}}}{1 - e^{-\frac{\Theta_{\text{vib}}}{T}}} \right)} = 1 - e^{-\frac{\Theta_{\text{vib}}}{T}}$$

$$\begin{aligned} P_{\text{excited}} &= 1 - P_0 \\ &= 1 - \left( 1 - e^{-\frac{\Theta_{\text{vib}}}{T}} \right) = e^{-\frac{\Theta_{\text{vib}}}{T}} \end{aligned}$$

**Vibrational Temperature**  $\Theta_{\text{vib}} = \frac{h\nu}{k_B}$

- Temperature needed to achieve substantial excited state populations.
- If  $T < \Theta_{\text{vib}}$ , the system is predominantly in the ground state
- If  $T \gtrsim \Theta_{\text{vib}}$ , there is substantial population in excited states

## Example: Vibrational Temperature

At 298 K, are molecules mostly in the ground or excited vibrational state?

Molecule	$\Theta_{\text{vib}}$ (K)
H <sub>2</sub>	6215
Cl <sub>2</sub>	80
I <sub>2</sub>	308

Excited-state probability:

$$P_{\text{excited}} = e^{-\Theta_{\text{vib}}/T}$$

Calculations:

$$P_{\text{excited}}(\text{H}_2) = e^{-6215/298} \approx 1.01 \times 10^{-19} \quad \text{Mostly ground state}$$

$$P_{\text{excited}}(\text{Cl}_2) = e^{-805/298} \approx 6.82 \times 10^{-2} \quad 7\% \text{ in excited state}$$

$$P_{\text{excited}}(\text{I}_2) = e^{-308/298} \approx 3.58 \times 10^{-1} \quad 36\% \text{ in excited state}$$

# Rotational Partition Function

Rigid Rotor:

$$\varepsilon_J = \frac{\hbar^2 J(J+1)}{2I} \quad J = 0, 1, 2, \dots$$

Each energy level has degeneracy:

$$g_J = 2J + 1 \quad m_J = -J, -J + 1, \dots, J - 1, J$$

Define rotational temperature:

$$\Theta_{\text{rot}} = \frac{\hbar^2}{2Ik_B}$$

Partition function:

$$q_{\text{rot}} = \sum_i g e^{-\beta E_i} = \sum_{J=0}^{\infty} (2J + 1) e^{-\frac{\Theta_{\text{rot}}(J(J+1))}{T}}$$

# Rotational Partition Function

$$q_{\text{rot}} = \sum_{J=0}^{\infty} (2J + 1) e^{-\Theta_{\text{rot}} J(J+1)/T}$$

Substitute  $u = J(J + 1)$ , so  $du = (2J + 1) dJ$ :

$$= \int_0^{\infty} e^{-u\Theta_{\text{rot}}/T} du$$

Using the integral  $\int_0^{\infty} e^{-au} du = \frac{1}{a}$  for  $a > 0$ :

$$= \frac{1}{\Theta_{\text{rot}}/T}$$
$$q_{\text{rot}} = \frac{T}{\Theta_{\text{rot}}} \quad T \gg \Theta_{\text{rot}} \quad (\text{high } T \text{ limit})$$

## What is Rotational Temperature?

Consider the probability of finding a particle in an excited state:

$$P_{excited} = 1 - P_0$$

Ground State energy:

$$\varepsilon_0 = \frac{\hbar^2}{2I} \cdot 0(0 + 1) = 0$$

Probability of being in the ground state:

$$P_0 = \frac{g_0 e^{-\beta \varepsilon_0}}{q_{rot}} = \frac{1 \cdot e^0}{T/\Theta_{rot}} = \frac{\Theta_{rot}}{T}$$

# What is Rotational Temperature?

Probability of being in excited rotational states:

$$\begin{aligned}P_{excited} &= 1 - P_0 \\ &= 1 - \frac{\Theta_{rot}}{T}\end{aligned}$$

**Rotational Temperature**  $\Theta_{rot} = \frac{\hbar^2}{2Ik_B}$

- Temperature needed to achieve substantial excited state populations.
- If  $T < \Theta_{rot}$ , the system is predominantly in the ground state
- If  $T \gtrsim \Theta_{rot}$ , there is substantial population in excited states

## Example: Rotational Temperature

At 298 K, are molecules mostly in the ground or excited rotational state?

Molecule	$\Theta_{\text{rot}}$ (K)
H <sub>2</sub>	87.5
CO	2.77
I <sub>2</sub>	0.054

Excited-state probability:

$$P_{\text{excited}} = 1 - \frac{\Theta_{\text{rot}}}{T}$$

Calculations:

$$P_{\text{excited}}(\text{H}_2) = 1 - \frac{87.5}{298} \approx 0.706 \quad 71\% \text{ in excited states}$$

$$P_{\text{excited}}(\text{CO}) = 1 - \frac{2.77}{298} \approx 0.991 \quad 99\% \text{ in excited states}$$

$$P_{\text{excited}}(\text{I}_2) = 1 - \frac{0.054}{298} \approx 0.9998 \quad >99.9\% \text{ in excited states}$$

# Vibrational Energy Contribution

$$q_{\text{vib}} = \frac{1}{1 - e^{-\Theta_{\text{vib}}/T}}$$

$$\begin{aligned}\langle \epsilon_{\text{vib}} \rangle &= -\frac{1}{q_{\text{vib}}} \frac{\partial}{\partial \beta} \left( \frac{1}{1 - e^{-\beta \hbar \omega}} \right) \\ &= -\frac{1}{q_{\text{vib}}} \frac{\partial}{\partial \beta} \left( 1 - e^{-\beta \hbar \omega} \right)^{-1} \\ &= -\frac{1}{q_{\text{vib}}} (-1) \left( 1 - e^{-\beta \hbar \omega} \right)^{-2} \left( -(-\hbar \omega) e^{-\beta \hbar \omega} \right) \\ &= (1 - e^{-\beta \hbar \omega}) \cdot \frac{\hbar \omega e^{-\beta \hbar \omega}}{(1 - e^{-\beta \hbar \omega})^2}\end{aligned}$$

$$\begin{aligned}\langle \epsilon_{\text{vib}} \rangle &= \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \\ &= \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \cdot \frac{e^{\beta \hbar \omega}}{e^{\beta \hbar \omega}} \\ &= \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1}\end{aligned}$$

$$e^x \approx 1 + x \text{ (for small } x\text{)}$$

$$e^{\beta \hbar \omega} \approx 1 + \beta \hbar \omega$$

$$\langle \epsilon_{\text{vib}} \rangle \approx \frac{\hbar \omega}{\beta \hbar \omega} = k_B T$$

$k_B T$  for each vibrational degree of freedom in diatomic, high T limit

## Rotational Energy Contribution

$$q_{\text{rot}} = \frac{T}{\Theta_{\text{rot}}}$$

$$\langle \varepsilon_{\text{rot}} \rangle = k_B T^2 \frac{1}{q_{\text{rot}}} \frac{\partial q_{\text{rot}}}{\partial T}$$

$$\frac{\partial q_{\text{rot}}}{\partial T} = \frac{\partial}{\partial T} \left( \frac{T}{\Theta_{\text{rot}}} \right) = \frac{1}{\Theta_{\text{rot}}}$$

$$\langle \varepsilon_{\text{rot}} \rangle = k_B T^2 \frac{1}{T/\Theta_{\text{rot}}} \cdot \frac{1}{\Theta_{\text{rot}}} = k_B T$$

$\frac{1}{2}k_B T$  per rotational degree of freedom in diatomic

# Summary: Molecular Partition Functions

## Summary Table

Mode	Partition function	Energy spacing	Population at 300 K
Trans	$q_{\text{trans}} = \left(\frac{2\pi m}{h^2\beta}\right)^{3/2} V$	Continuous ( $k_B T \gg \Delta\varepsilon$ )	More excited
Rot	$q_{\text{rot}} = \frac{T}{\Theta_{\text{rot}}}$	Small ( $\Theta_{\text{rot}} \sim 1-10$ K)	Mostly excited
Vib	$q_{\text{vib}} = \frac{e^{-\Theta_{\text{vib}}/2T}}{1 - e^{-\Theta_{\text{vib}}/T}}$	Large ( $\Theta_{\text{vib}} \sim 300-6000$ K)	Mostly ground
Elec	$q_{\text{elec}} = \sum_i g_i e^{-\beta\varepsilon_i}$	Very large ( $k_B T \ll \Delta\varepsilon$ )	Ground state

$$P_{\text{vib,excited}} = e^{-\Theta_{\text{vib}}/T} \quad P_{\text{rot,excited}} = 1 - \frac{\Theta_{\text{rot}}}{T} \quad \Theta_{\text{vib}} = \frac{h\nu}{k_B} \quad \Theta_{\text{rot}} = \frac{\hbar^2}{2Ik_B}$$

$T \gg \Theta \implies$  excited states populated       $T \ll \Theta \implies$  ground state dominates